

Isogeometric Analysis for Maxwell's equation

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Section 1

Introduction

Maxwell's Equation

$$\frac{\partial \mathbf{B}}{\partial t} + \operatorname{curl} \mathbf{E} = 0, \quad (\text{Faraday's law})$$

$$\frac{\partial \mathbf{D}}{\partial t} - \operatorname{curl} \mathbf{H} = -\mathbf{J}, \quad (\text{Ampère's law})$$

$$\operatorname{div} \mathbf{D} = \rho, \quad (\text{electrical Gauss law})$$

$$\operatorname{div} \mathbf{B} = 0, \quad (\text{magnetic Gauss law}).$$

Material laws

$$\mathbf{D} = \epsilon \mathbf{E},$$

$$\mathbf{B} = \nu(\mathbf{H} + \mathbf{M}),$$

$$\mathbf{J} = \sigma \mathbf{E}.$$

Quantities

\mathbf{B} .. mag. field \mathbf{J} .. el. current

\mathbf{E} .. el. field ρ .. el. charge density

\mathbf{D} .. el. displacement ϵ .. el. permittivity

\mathbf{H} .. mag. induction ν .. mag. permeability

\mathbf{M} .. magnetization σ .. conductivity



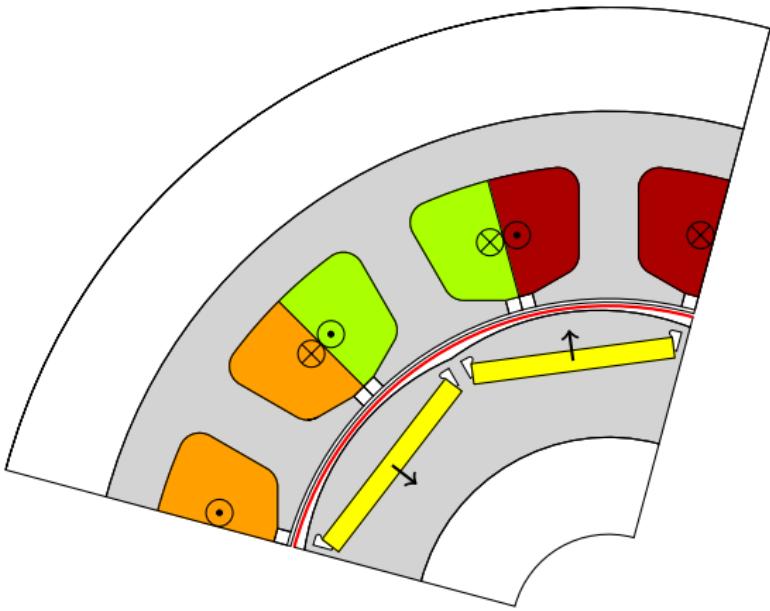


Figure: Motor sector

References I

-  A. Buffa, J. Rivas, G. Sangalli, and R. Vázquez.
Isogeometric discrete differential forms in three dimensions.
SIAM Journal on Numerical Analysis, 49(2):818–844, 2011.
-  A. Buffa, G. Sangalli, and R. Vázquez.
Isogeometric methods for computational electromagnetics:
B-spline and t-spline discretizations.
J. Comput. Phys., 257:1291–1320, January 2014.
-  A. Buffa, G. Sangalli, and R. Vázquez.
Isogeometric analysis in electromagnetics: B-splines
approximation.
Computer Methods in Applied Mechanics and Engineering,
199(17–20):1143 – 1152, 2010.

References II

-  Jie Li, Daniel Dault, Beibei Liu, Yiyang Tong, and Balasubramaniam Shanker.
Subdivision based isogeometric analysis technique for electric field integral equations for simply connected structures.
Journal of Computational Physics, 319:145 – 162, 2016.
-  Ahmed Ratnani and Eric Sonnendrücker.
An arbitrary high-order spline finite element solver for the time domain maxwell equations.
Journal of Scientific Computing, 51(1):87–106, 2012.
-  R. Vazquez and A. Buffa.
Isogeometric analysis for electromagnetic problems.
IEEE Transactions on Magnetics, 46(8):3305–3308, Aug 2010.

Eigenvalue Problem

Find $\omega \in \mathbb{R}$ and $u \in H_0(\text{curl}; \Omega)$ such that

$$\int_{\Omega} \mu^{-1} \operatorname{curl} u \operatorname{curl} v = \omega^2 \int_{\Omega} \epsilon u \cdot v \quad \forall v \in H_0(\text{curl}; \Omega).$$

Source Problem

Find $u \in H_0(\text{curl}; \Omega)$ such that

$$\int_{\Omega} \mu^{-1} \operatorname{curl} u \cdot \operatorname{curl} v - \omega^2 \int_{\Omega} \epsilon u \cdot v = \int_{\Omega} f \cdot v \quad \forall v \in H_0(\text{curl}; \Omega).$$

physical Domain Ω , and Parametrization \mathbf{F}

- $\Omega \subset \mathbb{R}^3$: bounded, simply connected Lipschitz domain with
 $\partial\Omega$: connected boundary
 $\mathbf{F} : \widehat{\Omega} \rightarrow \Omega$: continuously differentiable geometrical mapping
with continuously differentiable inverse

Sobolev spaces

$$\mathbf{H}(\mathbf{curl}; \Omega) := \{\mathbf{v} \in \mathbf{L}^2(\Omega) \mid \mathbf{curl}(\mathbf{v}) \in \mathbf{L}^2(\Omega)\}$$

$$\mathbf{H}(\text{div}; \Omega) := \{\mathbf{v} \in \mathbf{L}^2(\Omega) \mid \text{div}(\mathbf{v}) \in \mathbf{L}^2(\Omega)\}$$

De Rham Complex

$$\mathbb{R} \longrightarrow H^1(\widehat{\Omega}) \xrightarrow{\widehat{\text{grad}}} \mathbf{H}(\mathbf{curl}; \widehat{\Omega}) \xrightarrow{\widehat{\text{curl}}} \mathbf{H}(\text{div}; \widehat{\Omega}) \xrightarrow{\widehat{\text{div}}} L^2(\widehat{\Omega}) \longrightarrow 0$$

$$\mathbb{R} \longrightarrow H^1(\Omega) \xrightarrow{\text{grad}} \mathbf{H}(\mathbf{curl}; \Omega) \xrightarrow{\text{curl}} \mathbf{H}(\text{div}; \Omega) \xrightarrow{\text{div}} L^2(\Omega) \longrightarrow 0$$

Exact for Ω (and $\widehat{\Omega}$) simply connected.

Pullback operators

$$\begin{aligned}\iota^0(\phi) &:= \phi \circ \mathbf{F}, & \phi &\in H^1(\Omega) \\ \iota^1(\mathbf{u}) &:= (D\mathbf{F})^T(\mathbf{u} \circ \mathbf{F}), & \mathbf{u} &\in \mathbf{H}(\mathbf{curl}; \Omega) \\ \iota^2(\mathbf{v}) &:= \det(D\mathbf{F})(D\mathbf{F})^{-1}(\mathbf{v} \circ \mathbf{F}), & \mathbf{v} &\in \mathbf{H}(\operatorname{div}; \Omega) \\ \iota^3(\varphi) &:= \det(D\mathbf{F})(\varphi \circ \mathbf{F}), & \varphi &\in L^2(\Omega)\end{aligned}$$

De Rham Complex

$$\begin{array}{ccccccc} \mathbb{R} & \longrightarrow & H^1(\widehat{\Omega}) & \xrightarrow{\widehat{\text{grad}}} & \mathbf{H}(\mathbf{curl}; \widehat{\Omega}) & \xrightarrow{\widehat{\text{curl}}} & \mathbf{H}(\text{div}; \widehat{\Omega}) & \xrightarrow{\widehat{\text{div}}} & L^2(\widehat{\Omega}) \longrightarrow 0 \\ & & \iota^0 \uparrow & & \iota^1 \uparrow & & \iota^2 \uparrow & & \iota^3 \uparrow \\ \mathbb{R} & \longrightarrow & H^1(\Omega) & \xrightarrow{\text{grad}} & \mathbf{H}(\mathbf{curl}; \Omega) & \xrightarrow{\text{curl}} & \mathbf{H}(\text{div}; \Omega) & \xrightarrow{\text{div}} & L^2(\Omega) \longrightarrow 0 \end{array}$$

Section 2

B-Splines

Notation

$$\Sigma := \{\xi_1, \dots, \xi_{n+p+1}\}$$
 p-open knot vector
$$\Sigma' := \{\xi_2, \dots, \xi_{n+p}\}$$
$$B_{i,p}(\xi)$$

B-spline functions

$$S_p(\Sigma) := \text{span}\{B_{i,p}, i = 1, \dots, n\}$$

Anchors and Greville Sites

Anchors	Greville sites
$\xi^A := \begin{cases} \xi_{i+\frac{p+1}{2}} & p \text{ odd,} \\ \frac{\xi_{i+\frac{p}{2}} + \xi_{i+\frac{p}{2}+1}}{2} & p \text{ even} \end{cases}$	$\gamma^A = \frac{\xi_{i+1} + \dots + \xi_{i+p}}{p}$

Multivariate Case

$$\Sigma_i := \{\xi_{i,1}, \dots, \xi_{i,n_i+p_i+1}\}$$

$$\mathcal{A}_{p_1, \dots, p_d}(\Sigma_1, \dots, \Sigma_d) := \mathcal{A}_{p_1}(\Sigma_1) \times \dots \times \mathcal{A}_{p_d}(\Sigma_d)$$

$$B_{p_1, \dots, p_d}^A(\xi) = B_{p_1}^{A_1}(\xi_1) \dots B_{p_d}^{A_d}(\xi_d)$$

$$S_{p_1, \dots, p_d}(\Sigma_1, \dots, \Sigma_d) := \text{span}\{B_{p_1, \dots, p_d}^A\}$$

Spline spaces will be high-order extensions of classical low order Nédélec hexahedral finite elements

Discrete Spaces

$$\widehat{X}_h^0 := S_{p_1, p_2, p_3}(\Sigma_{1,2,3}),$$

$$\widehat{X}_h^1 := S_{p_1-1, p_2, p_3}(\Sigma'_1, \Sigma_{2,3}) \times S_{p_1, p_2-1, p_3}(\Sigma_1, \Sigma'_2, \Sigma_3) \times S_{p_1, p_2, p_3-1}(\Sigma_{1,2}, \Sigma'_3),$$

$$\widehat{X}_h^2 := S_{p_1, p_{2,3}-1}(\Sigma_1, \Sigma'_{2,3}) \times S_{p_1-1, p_2, p_3-1}(\Sigma'_1, \Sigma_2, \Sigma'_3) \times S_{p_{1,2}-1, p_3}(\Sigma'_{1,2}, \Sigma_3),$$

$$\widehat{X}_h^3 := S_{p_1-1, p_2-1, p_3-1}(\Sigma'_1, \Sigma'_2, \Sigma'_3).$$

Discrete De Rham complex

$$\mathbb{R} \longrightarrow \widehat{X}_h^0 \xrightarrow{\widehat{\text{grad}}} \widehat{X}_h^1 \xrightarrow{\widehat{\text{curl}}} \widehat{X}_h^2 \xrightarrow{\widehat{\text{div}}} \widehat{X}_h^3 \longrightarrow 0$$

This sequence is exact.

Spaces on the Physical domain

Discrete De Rham complex

$$\begin{array}{cccc} \mathbb{R} \longrightarrow \widehat{X}_h^0 & \xrightarrow{\widehat{\text{grad}}} & \widehat{X}_h^1 & \xrightarrow{\widehat{\text{curl}}} \quad \widehat{X}_h^2 & \xrightarrow{\widehat{\text{div}}} \quad \widehat{X}_h^3 \longrightarrow 0 \\ \iota^0 \uparrow & & \iota^1 \uparrow & & \iota^2 \uparrow & & \iota^3 \uparrow \\ \mathbb{R} \longrightarrow X_h^0 & \xrightarrow{\text{grad}} & X_h^1 & \xrightarrow{\text{curl}} & X_h^2 & \xrightarrow{\text{div}} & X_h^3 \longrightarrow 0 \end{array}$$

Discrete Spaces

$$\begin{aligned} X_h^0 &:= \{\phi : \iota^0(\phi) \in \widehat{X}_h^0\}, \\ X_h^1 &:= \{\mathbf{u} : \iota^1(\mathbf{u}) \in \widehat{X}_h^1\}, \\ X_h^2 &:= \{\mathbf{v} : \iota^2(\mathbf{v}) \in \widehat{X}_h^2\}, \\ X_h^3 &:= \{\varphi : \iota^3(\varphi) \in \widehat{X}_h^3\}. \end{aligned}$$

Projectors onto the B-spline space

$$\widehat{\Pi}^{p_1, p_2, p_3} \phi := \sum_{i_1=2, i_2=2, i_3=2}^{n_1-1, n_2-1, n_3-1} (\lambda_{i_1, i_2, i_3} \phi) B_{i_1, i_2, i_3}$$

where λ_i^p are the dual basis functionals in each variable:

$$\lambda_i^p B_j^p = \delta_{ij}$$

Spline preserving property

$$\widehat{\Pi}^0 \widehat{\phi}_h := \widehat{\phi}_h, \quad \forall \widehat{\phi}_h \in \widehat{X}_h^0$$

$$\widehat{\Pi}^1 \widehat{\mathbf{u}}_h := \widehat{\mathbf{u}}_h, \quad \forall \widehat{\mathbf{u}}_h \in \widehat{X}_h^1$$

$$\widehat{\Pi}^2 \widehat{\mathbf{v}}_h := \widehat{\mathbf{v}}_h, \quad \forall \widehat{\mathbf{v}}_h \in \widehat{X}_h^2$$

$$\widehat{\Pi}^3 \widehat{\varphi}_h := \widehat{\varphi}_h, \quad \forall \widehat{\varphi}_h \in \widehat{X}_h^3$$



Projectors onto the B-spline space

$$\begin{array}{ccccccc} \mathbb{R} & \longrightarrow & H^1(\widehat{\Omega}) & \xrightarrow{\widehat{\text{grad}}} & \mathbf{H}(\mathbf{curl}; \widehat{\Omega}) & \xrightarrow{\widehat{\text{curl}}} & \mathbf{H}(\text{div}; \widehat{\Omega}) & \xrightarrow{\widehat{\text{div}}} & L^2(\widehat{\Omega}) \longrightarrow 0 \\ & & \widehat{\Pi}_0^0 \downarrow & & \widehat{\Pi}_0^1 \downarrow & & \widehat{\Pi}_0^2 \downarrow & & \widehat{\Pi}_0^3 \downarrow \\ \mathbb{R} & \longrightarrow & \widehat{X}_h^0 & \xrightarrow{\widehat{\text{grad}}} & \widehat{X}_h^1 & \xrightarrow{\widehat{\text{curl}}} & \widehat{X}_h^2 & \xrightarrow{\widehat{\text{div}}} & \widehat{X}_h^3 \longrightarrow 0 \end{array}$$

The same holds also for the physical domain.

Approximation estimate

Under the assumption $\gamma_d \geq \alpha_d$, $d = 1, 2, 3$, the following estimates hold:

$$\|\phi - \Pi^0 \phi\|_{H^l(\Omega)} \leq Ch^{s-l} \|\phi\|_{H^s(\Omega)} \quad \forall \phi \in H^1 \cap H^s(\Omega),$$

$$0 \leq l \leq s \leq p+1, l \leq \alpha+1$$

$$\|\mathbf{u} - \Pi^1 \mathbf{u}\|_{H^l(\Omega)} \leq Ch^{s-l} \|\mathbf{u}\|_{H^s(\Omega)} \quad \forall \mathbf{u} \in \mathbf{H}(\mathbf{curl}; \Omega) \cap \mathbf{H}^s(\Omega),$$

$$0 \leq l \leq s \leq p, l \leq \alpha$$

$$\|\mathbf{v} - \Pi^2 \mathbf{v}\|_{H^l(\Omega)} \leq Ch^{s-l} \|\mathbf{v}\|_{H^s(\Omega)} \quad \forall \mathbf{v} \in \mathbf{H}(\text{div}; \Omega) \cap \mathbf{H}^s(\Omega),$$

$$0 \leq l \leq s \leq p, l \leq \alpha$$

$$\|\varphi - \Pi^3 \varphi\|_{H^l(\Omega)} \leq Ch^{s-l} \|\varphi\|_{H^s(\Omega)} \quad \forall \varphi \in L^2(\Omega) \cap H^s(\Omega),$$

$$0 \leq l \leq s \leq p, l \leq \alpha.$$

Approximation estimate (Energy Norm)

The following inequalities hold for $0 \leq l \leq s \leq p, l \leq \alpha$:

$$\|\phi - \Pi^0 \phi\|_{H^{l+1}(\Omega)} \leq Ch^{s-l} \|\phi\|_{H^{s+1}(\Omega)}$$
$$\forall \phi \in H^{s+1}(\Omega),$$

$$\|\mathbf{u} - \Pi^1 \mathbf{u}\|_{\mathbf{H}^l(\mathbf{curl}; \Omega)} \leq Ch^{s-l} \|\mathbf{u}\|_{\mathbf{H}^s(\mathbf{curl}; \Omega)}$$
$$\forall \mathbf{u} \in \mathbf{H}^s(\mathbf{curl}; \Omega),$$

$$\|\mathbf{v} - \Pi^2 \mathbf{v}\|_{\mathbf{H}^l(\mathbf{div}; \Omega)} \leq Ch^{s-l} \|\mathbf{v}\|_{\mathbf{H}^s(\mathbf{div}; \Omega)}$$
$$\forall \mathbf{v} \in \mathbf{H}^s(\mathbf{div}; \Omega),$$

$$\|\varphi - \Pi^3 \varphi\|_{H^l(\Omega)} \leq Ch^{s-l} \|\varphi\|_{H^s(\Omega)}$$
$$\forall \varphi \in H^s(\Omega).$$

Conformity across the interface

If the De-Rham complex is fulfilled on each patch:

- trace continuity on X_h^0
- tangential trace continuity on X_h^1
- normal trace continuity on X_h^2
- no continuity on X_h^3

Geometrical Conformity

On each non-empty patch interface Γ the spaces X_{h,k_1}^0 and X_{h,k_2}^0 coincide, as the corresponding bases do.

Section 3

T-Splines

Notation

$$\Sigma_i := \{\xi_{i,1}, \dots, \xi_{i,n_i+p_i+1}\}$$

$$\mathcal{A}_{p_1, \dots, p_d}(\Sigma_1, \dots, \Sigma_d) := \mathcal{A}_{p_1}(\Sigma_1) \times \dots \times \mathcal{A}_{p_d}(\Sigma_d)$$

$$B_{p_1, \dots, p_d}^A(\xi) = B[\Sigma_1^A](\xi_1) \dots B[\Sigma_d^A](\xi_d)$$

$$T_{p_1, \dots, p_d}(\mathcal{M}) := \text{span}\{B_{p_1, \dots, p_d}^A : A \in \mathcal{A}_{p_1, \dots, p_d}(\mathcal{M})\}$$

2D De Rham Complex

We have to modify the mesh not only at the boundary but also at the T-junctions.

$$\widehat{Y}_h^0 := T_{p,p}(\mathcal{M}^0)$$

$$\widehat{Y}_h^1 := T_{p-1,p}(\mathcal{M}_1^1) \times T_{p,p-1}(\mathcal{M}_2^1)$$

$$\widehat{Y}_h^{1*} := T_{p,p-1}(\mathcal{M}_2^1) \times T_{p-1,p}(\mathcal{M}_1^1)$$

$$\widehat{Y}_h^2 := T_{p-1,p-1}(\mathcal{M}^2)$$

De Rham Complex

$$\mathbb{R} \longrightarrow \widehat{Y}_h^0 \xrightarrow{\widehat{\text{grad}}} \widehat{Y}_h^1 \xrightarrow{\widehat{\text{curl}}} \widehat{Y}_h^2 \longrightarrow 0$$

$$\mathbb{R} \longrightarrow \widehat{Y}_h^0 \xrightarrow{\widehat{\text{curl}}} \widehat{Y}_h^{1*} \xrightarrow{\widehat{\text{div}}} \widehat{Y}_h^2 \longrightarrow 0$$

3D De Rham Complex

$$\widehat{X}_h^0 := \widehat{Y}_h^0 \times S_p(\Sigma)$$

$$\widehat{X}_h^1 := [\widehat{Y}_h^1 \times S_p(\Sigma)] \times [\widehat{Y}_h^0 \times S_{p-1}(\Sigma')]$$

$$\widehat{X}_h^2 := [\widehat{Y}_h^{1*} \times S_{p-1}(\Sigma')] \times [\widehat{Y}_h^2 \times S_p(\Sigma)]$$

$$\widehat{Y}_h^3 := \widehat{Y}_h^2 \times S_{p-1}(\Sigma')$$

De Rham Complex

$$\mathbb{R} \longrightarrow \widehat{X}_h^0 \xrightarrow{\widehat{\text{grad}}} \widehat{X}_h^1 \xrightarrow{\widehat{\text{curl}}} \widehat{X}_h^2 \xrightarrow{\widehat{\text{div}}} \widehat{X}_h^3 \longrightarrow 0$$

Section 4

Examples

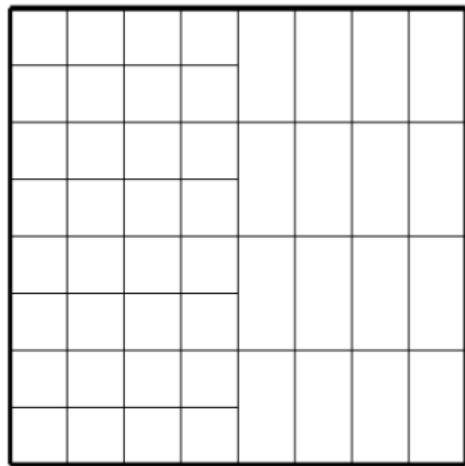
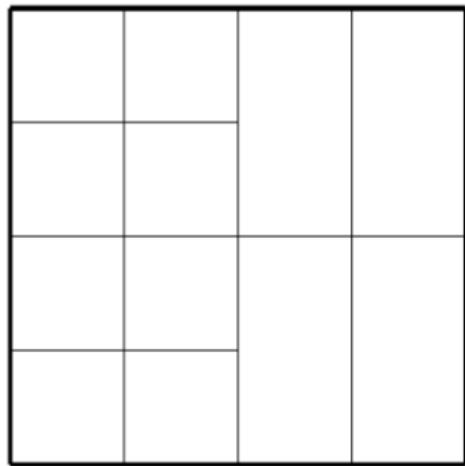
Eigenvalue Problem

Find $\omega \in \mathbb{R}$ and $u \in H_0(\text{curl}; \Omega)$ such that

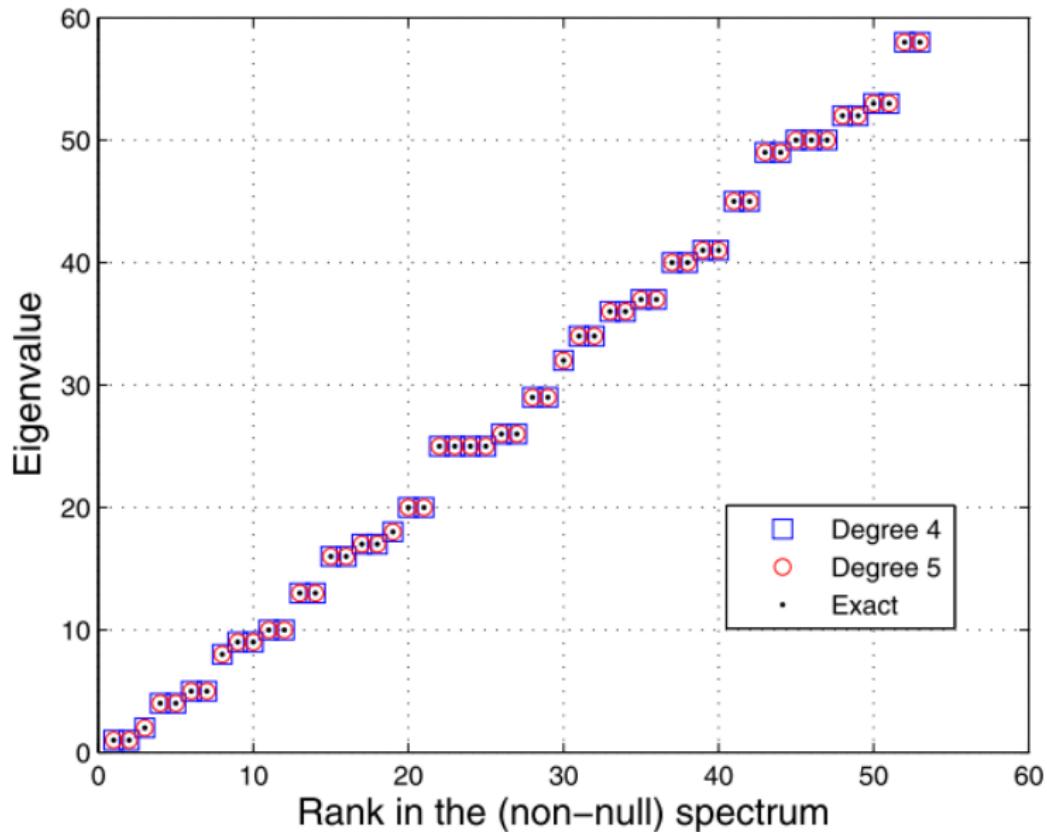
$$\int_{\Omega} \mu^{-1} \operatorname{curl} u \operatorname{curl} v = \omega^2 \int_{\Omega} \epsilon u \cdot v \quad \forall v \in H_0(\text{curl}; \Omega).$$

Aim is to show that there are no spurious modes with T-splines.

Results: Square

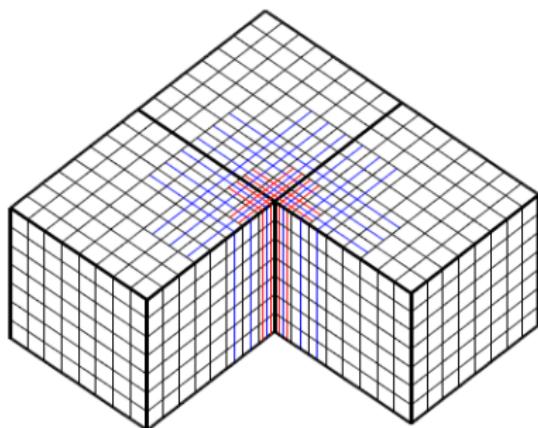
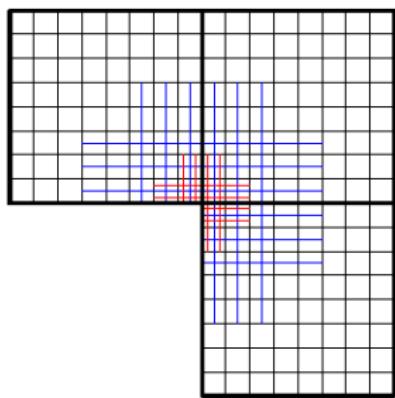


Results: Square

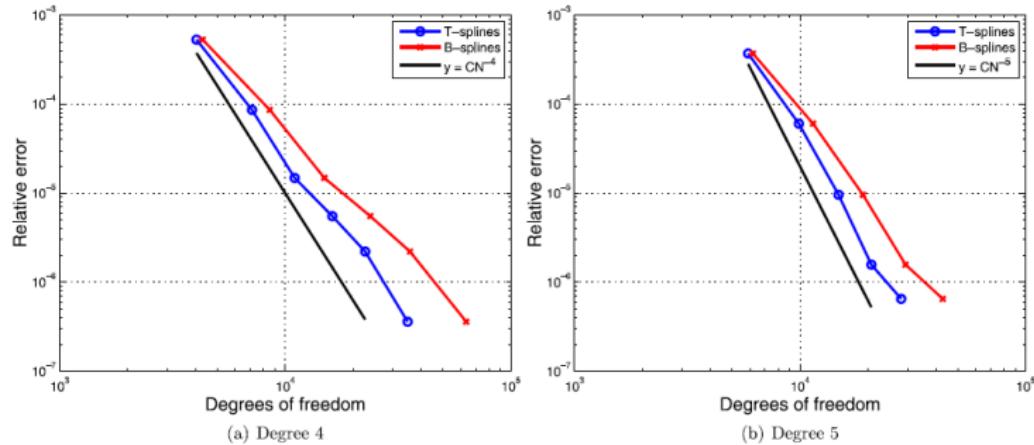


Results: L-Shape Domain 3d

With suitable refinement due to the reentrant edge:



Results: L-Shape Domain 3d



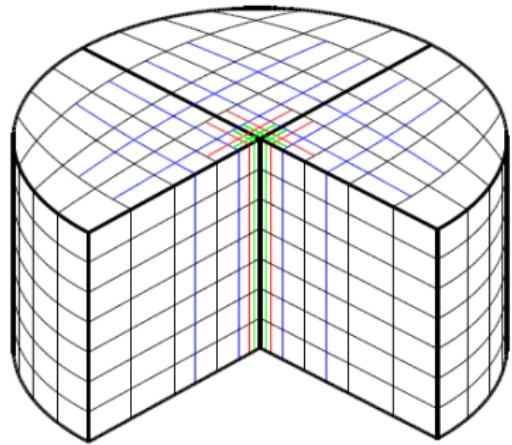
The method is free of spurious modes.

Source Problem

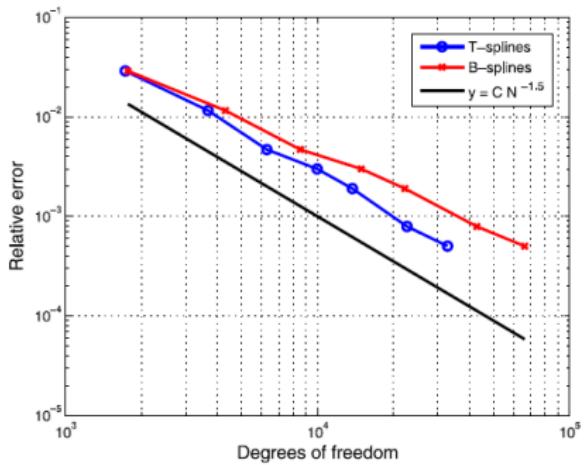
Find $u \in H_0(\text{curl}; \Omega)$ such that

$$\int_{\Omega} \mu^{-1} \operatorname{curl} u \cdot \operatorname{curl} v - \omega^2 \int_{\Omega} \epsilon u \cdot v = \int_{\Omega} f \cdot v \quad \forall v \in H_0(\text{curl}; \Omega),$$

Results: Three quarters of a cylinder



T- and B-spline of degree 3.



Thank you for your attention!